# Non-smooth pitchfork bifurcations in a quasi-periodically forced piecewise-linear map 

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#### Abstract

We study a class of one-dimensional family of quasiperiodically forced maps $F_{a, b}(x, \theta)=$ $\left(f_{a, b}(x, \theta), \theta+\omega\right)$, where $x$ is real, $\theta$ is an angle, and $\omega$ is an irrational frequency, such that $f_{a, b}(x, \theta)$ is a real piecewise linear map with respect to $x$ of certain kind. The family depends on two real parameters, $a>0$ and $b>0$. For this family, we prove the existence of nonsmooth pitchfork bifurcations. For $a<1$ and any $b$ there is only a continuous invariant curve. For $a>1$ there exists a smooth map $b=b_{0}(a)$ such that: a) For $b<b_{0}(a), f_{a, b}$ has two continuous attracting invariant curves and one continuous repelling one; b) For $b=b_{0}(a)$ it has one continuous repelling invariant curve and two semicontinuous (non-continuous) attracting invariant curves that intersect the unstable one in a zero-Lebesgue measure set of angles; c) For $b>b_{0}(a)$ it has one continuous attracting invariant curve. The case $a=1$ is a degenerate case that is also discussed in the paper. It is interesting to note that this family is a simplified version of the smooth family $G_{a, b}(x, \theta)=(\arctan (a x)+b \sin (\theta), \theta+\omega)$ for which there is numerical evidence of a non-smooth pitchfork bifurcation. Finally, we also discuss the limit case when $a \rightarrow \infty$.


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## 1 Introduction

In this paper we study the dynamics of the quasiperiodically forced map $(\bar{x}, \bar{\theta})=F_{a, b}(x, \theta)$, where

$$
\left\{\begin{array}{l}
\bar{x}=h_{a}(x)+b \sin \theta,  \tag{1}\\
\bar{\theta}=\theta+\omega \bmod 2 \pi,
\end{array}\right.
$$

for $(x, \theta) \in \mathbb{R} \times \mathbb{T}^{1}, \omega \notin 2 \pi \mathbb{Q}$ and $h_{a}$ is the continuous piecewise linear map

$$
h_{a}(x)=\left\{\begin{array}{rll}
-\frac{\pi}{2} & \text { if } & x \leq-\frac{\pi}{2 a}  \tag{2}\\
a x & \text { if } & -\frac{\pi}{2 a}<x<\frac{\pi}{2 a} \\
\frac{\pi}{2} & \text { if } & x \geq \frac{\pi}{2 a}
\end{array}\right.
$$

This dynamical system depends on two parameters: $b$ is a real value and $a$ is a real value strictly larger than 0 .

The motivation to study this map comes from [JMAT18], that studies the map

$$
\left\{\begin{array}{l}
\bar{x}=\arctan (a x)+b \sin \theta  \tag{3}\\
\bar{\theta}=\theta+\omega \quad \bmod 2 \pi
\end{array}\right.
$$

which is a rescaled version of a map already studied in [Jäg03]. First, assume that $a>1$ and that $b=0$. In this case, it is clear that the map $x \mapsto \arctan (a x)$ has exactly three fixed points, that are seen in (3) as three (constant) invariant curves, one repelling and two attracting. When $b$ becomes different from zero (but small), the three invariant curves can be continued w.r.t. $b$ (for instance, using the Implicit Function Theorem) and they become three non-constant invariant curves, one repelling and two attracting. Note that (3) is invariant by the symmetry $S:(x, \theta) \mapsto(-x, \theta+\pi)$ and, therefore, if $\theta \mapsto \varphi(\theta)$ is an invariant curve of (3),

$$
\varphi(\theta+\omega)=\arctan (a \varphi(\theta))+b \sin \theta
$$

then $\theta \mapsto-\varphi(\theta+\pi)$ is also an invariant curve of (3). As the system cannot have more than three invariant curves (see also [Jäg03]), the repelling curve is self-symmetric and the attracting curves are one the image of the other by the symmetry $S$. If $a$ is small enough (but larger than 1 ), when the value of $b$ increases the three invariant curves meet in a pitchfork bifurcation and, after that, only one self-symmetric attracting curve exists ([Jäg03, JMAT18]). The situation seems to be different when $a$ is large. In this case, when $b$ is increased and the three invariant curve approach, they start to wrinkle and they seem (numerically) to become a strange set when they merge for a critical value $b^{*}$. After the merging, if we increase again the value of $b$, the strange set becomes a smooth (but very wrinkled) attracting curve, see [JMAT18].

The map $F_{a, b}$ given by (1) is a "simplified" version of (3) that aims to approximate the main dynamical features of (1), specially for large values of $a$. We also note that (1) has the same symmetry $S:(x, \theta) \mapsto(-x, \theta+\pi)$ as (3).

Figure 1 shows the attracting sets of (1) for several values of the parameters. For these plots we have chosen $\omega$ to be the golden mean $(\omega=\pi(\sqrt{5}-1))$ and $a=4$. The first plot $(b=1.8)$ shows two attracting curves that are 'mirror' images under the symmetry $S$. As we will show


Figure 1: Attracting sets of (1) for different values of the parameters. The horizontal axis is $\theta$ and the vertical is $x$. The frequency $\omega$ is the golden mean.
later, there is a repelling self-symmetric invariant curve between these two. When the parameter $b$ is increased, the two attracting curves start to wrinkle and to approach the repelling curve and it seems that they merge into a single self-symmetric attracting curve as it happens in the pitchfork bifurcation but, in this case, the merging process seems to happen through a strange attracting set.

In this paper we prove that, for each $\omega \notin 2 \pi \mathbb{Q}$ and for $a>1$, there exists a critical value $b^{*}=b^{*}(a, \omega)$, such that when $0<b<b^{*}$ the map $F_{a, b}$ has three continuous invariant curves, two attracting (that are mirror images under $S$ ) and one repelling which is self-symmetric. When $b=b^{*}$ the map has only one continuous curve (which is repelling and self-symmetric), plus two semicontinuous invariant curves that are mirror images under $S$. Each semicontinuous curve intersects the continuous repelling curve in a dense set of points, and the $\theta$ coordinates of these points have zero measure in $[0,2 \pi]$. The semicontinuous curve above the repelling one is upper semicontinuous and the one below the repelling curve is lower semicontinuous. Finally, when $b>b^{*}$, we show that there exists a unique self-symmetric invariant curve. In other words, we show the existence of a non-smooth pitchfork bifurcation, in the sense that for $b<b^{*}$ there are three continuous invariant graphs, for $b>b^{*}$ there is only one continuous graph and for the critical value $b=b^{*}$ there exists a strange invariant set that can be described as follows: the upper boundary is given by an attracting upper semicontinuous graph and the lower boundary is given by an attracting lower semicontinuous graph. Moreover, the set also contains a continuous repelling curve, and the two semicontinuous attracting graphs intersect the continuous repelling curve on a dense set of values of $\theta$. This set is also invariant by the symmetry $S$.

The case $0<a \leq 1$ is also considered. We show that, for $a<1$ there is only one invariant curve which is attracting. The case $a=1$ is a bit more involved due to the neutral character of the map around the origin. We show that for $b<b^{*}(1, \omega)$ there is a one parametric family of invariant curves (parametrized over a nontrivial closed interval) and, for $b \geq b^{*}(1, \omega)$, there exists


Figure 2: The parameter space $(a, b)$ for the model (1). The blue zone corresponds to a single attracting self symmetric invariant curve, the green zone to three invariant curves, one repelling and self symmetric and two attracting curves that are mirror images under $S$. The red curve corresponds to SNAs. Crossing the red line when $a>1$ results in a non smooth pitchfork bifurcation.
a unique attracting invariant curve. All this is summarized in Figure 2 (see also Section 2.4 for another discussion of this figure).

A limit case is when $a \rightarrow+\infty$. To study this case we introduce the map

$$
h(x)=\left\{\begin{array}{rll}
-\frac{\pi}{2} & \text { if } & x<0  \tag{4}\\
0 & \text { if } & x=0 \\
\frac{\pi}{2} & \text { if } & x>0
\end{array}\right.
$$

which is the pointwise limit of $h_{a}(x)$ when $a \rightarrow \infty$ (and also the pointwise limit of the $\arctan (a x)$ function that appears in (3)). Figure 3 shows the attracting sets for different values of $b$. This case is simpler, but it still has some similarities with the previous one. Here, the critical value for $b$ is $\pi / 2$ and, as we will see, it is the limit of the previous value $b^{*}(a, \omega)$ when a tends to infinity. When $h_{a}=h$ and $b<\pi / 2$ the map (1) has two continuous invariant curves,

$$
\varphi_{0}(\theta)=\frac{\pi}{2}+b \sin (\theta-\omega), \quad \gamma_{0}(\theta)=-\frac{\pi}{2}+b \sin (\theta-\omega)
$$

Both are, in fact, superattracting since all initial conditions land on one of these curves in a finite number of iterates. As before, the curves are mirror images by the symmetry $S$. When $b$ reaches $\pi / 2$ these curves touch $x=0$ (at $\theta=3 \pi / 2+\omega$ for $\varphi_{0}$ and $\theta=\pi / 2+\omega$ for $\gamma_{0}$ ). There are two different situations, depending on the sign of $\sin (\pi / 2+\omega)$. If this sign is positive, there are two disjoint attracting curves that are discontinuous on one point, and if it is negative (this is the case when $\omega$ is the golden mean), the attracting set is the union of two attracting curves, with a dense set of discontinuities. It is remarkable that, in both cases, none of the attracting curves is invariant. Finally, when $b>\pi / 2$, there is only one invariant curve which is attracting


Figure 3: Attracting sets of (4) for different values of the parameters. The horizontal axis is $\theta$ and the vertical is $x$. The frequency $\omega$ is the golden mean.
and has a finite number of discontinuities. The number of discontinuities goes to infinity when $b$ tends to $\pi / 2$ from above.

Finally, Figure 4 shows the attracting sets of the different maps corresponding to (3), (1) and (4). Each column refers to one of the maps, the first row corresponds to $b=1$ and the second row to $b=10$. We have used the value $a=10$ to show the strong similarities between these maps when $a$ is large. This similarity is one of the motivations to study the piecewise linear map of this paper as an intermediate step to understand (3).

The existence of SNAs in quasiperiodically forced skew-product is known since the examples of [GOPY84] and [Kel96] (see also [JNOT07, LYL ${ }^{+}$20] and references therein). It is worth noting that numerical methods are quite limited to study SNAs, as pointed out in [JT08]. Studies of non-smooth saddle-node bifurcations in quite general situations can be found in [NDJS11, Fuh16, FGJ18]. A first investigation of non-smooth pitchfork bifurcations in a concrete model is contained in [Gle04].

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Figure 4: The first, second and third columns contain the attracting sets for the maps (3), (1) and the discontinuous map given by (4), respectively, for different values of the parameters. The horizontal axis is $\theta$ and the vertical is $x$. The frequency $\omega$ is the golden mean.

## 2 A piecewise linear quasiperiodically forced system

In this section we focus on the invariant curves of (1) with $h_{a}$ defined as in (2), with $\omega \notin 2 \pi \mathbb{Q}$. We consider $a>0$ but, to simplify the reading, the presentation has been split in the cases $a>1$ (Sections 2.1 and 2.2) and $a \geq 1$ (Section 2.3). As we will see, the values $a=1, b=b^{*}$ ( $b^{*}$ is defined below) is a critical point in the parameter space ( $a, b$ ) and it will be discussed in Section 2.4.

Now, let us assume that $a \neq 1$ (and $a>0$ ). Suppose that this map has an invariant curve $\eta$ such that $\eta(\theta) \in[-\pi /(2 a), \pi /(2 a)]$ for all $\theta \in \mathbb{T}^{1}$. Imposing the invariant condition,

$$
\begin{equation*}
\eta(\theta+\omega)=h_{a}(\eta(\theta))+b \sin \theta=a \eta(\theta)+b \sin \theta, \tag{5}
\end{equation*}
$$

we obtain a closed expression for $\eta$,

$$
\begin{equation*}
\eta(\theta)=-\frac{b \sin \omega}{(\cos \omega-a)^{2}+\sin ^{2} \omega} \cos \theta+\frac{b(\cos \omega-a)}{(\cos \omega-a)^{2}+\sin ^{2} \omega} \sin \theta \tag{6}
\end{equation*}
$$

The condition $\eta(\theta) \in[-\pi /(2 a), \pi /(2 a)]$ implies a restriction for the values of $a$ and $b$ : it is easy to see that the restriction is

$$
\begin{equation*}
|b| \leq b^{*}=b^{*}(a)=\frac{\pi}{2 a} \sqrt{1-2 a \cos \omega+a^{2}} . \tag{7}
\end{equation*}
$$

Moreover, there exist two angles $\theta_{ \pm}$such that $\eta^{\prime}\left(\theta_{ \pm}\right)=0$. We have that

$$
\cos \theta_{ \pm}= \pm \frac{\sin \omega}{\sqrt{1+a^{2}-2 a \cos \omega}}, \quad \sin \theta_{ \pm}=\mp \frac{\cos \omega-a}{\sqrt{1+a^{2}-2 a \cos \omega}}
$$

and

$$
\eta\left(\theta_{ \pm}\right)=\mp \frac{b}{\sqrt{1+a^{2}-2 a \cos \omega}} .
$$

In the limit case $b=b^{*}$, we have that $\eta\left(\theta_{ \pm}\right)=\mp \frac{\pi}{2 a}$ and $\eta^{\prime}\left(\theta_{ \pm}\right)=0$. This implies this curve is no longer invariant if $|b|>b^{*}$. Note that, for $a=1$ and $b<b^{*}, \eta$ is still an invariant curve but it is not the only invariant curve, and it is unique again for $a=1$ and $b=b^{*}$. The details are discussed in Section 2.3.

For any $a>0$, it follows that the image of the set $\left\{(x, \theta)\right.$ such that $\left.x \geq \pi /(2 a), \theta \in \mathbb{T}^{1}\right\}$ is the curve $\varphi_{0}(\theta)=\pi / 2+b \sin (\theta-\omega)$.
Lemma 1. If $a>0$ and $0 \leq b \leq b^{*}$, we have that

$$
\min _{\theta \in \mathbb{T}}\left(\varphi_{0}(\theta)-\eta(\theta)\right)=\frac{\pi}{2}-\frac{a b}{\sqrt{1+a^{2}-2 a \cos \omega}}=a \min _{\theta \in \mathbb{T}}\left(\frac{\pi}{2 a}-\eta(\theta)\right) .
$$

Moreover, the first minimum is attained at the single point $\theta_{0}=\theta_{-}+\omega$ and the second minimum is attained at the single point $\theta_{-}=\theta_{0}-\omega$.

Proof: By taking derivatives, it is easy to see that the angle $\theta_{0}$ for which we have the minimum distance satisfies:

$$
\cos \theta_{0}=\frac{\beta}{\sqrt{\alpha^{2}+\beta^{2}}}, \quad \sin \theta_{0}=-\frac{\alpha}{\sqrt{\alpha^{2}+\beta^{2}}}
$$

where

$$
\alpha=\frac{a b\left(1+a \cos \omega-2 \cos ^{2} \omega\right)}{1-2 a \cos \omega+a^{2}}, \quad \beta=\frac{a b \sin \omega(a-2 \cos \omega)}{1-2 a \cos \omega+a^{2}} .
$$

Moreover,

$$
\varphi_{0}\left(\theta_{0}\right)-\eta\left(\theta_{0}\right)=\frac{\pi}{2}-\sqrt{\alpha^{2}+\beta^{2}}
$$

and then

$$
\varphi_{0}\left(\theta_{0}\right)-\eta\left(\theta_{0}\right)=\frac{\pi}{2}-\frac{a b}{\sqrt{1+a^{2}-2 a \cos \omega}} .
$$

To see that $\theta_{0}=\theta_{-}+\omega$ we note that (5) implies

$$
\begin{equation*}
\varphi_{0}(\theta)-\eta(\theta)=a\left(\frac{\pi}{2 a}-\eta(\theta-\omega)\right), \quad \forall \theta \in \mathbb{T} . \tag{8}
\end{equation*}
$$

As the maximum of $\eta$ is attained at $\theta_{-}$, the minimum of $a(\pi /(2 a)-\eta(\theta))$ is attained at $\theta_{-}$, and using (8) we deduce that $\theta_{-}=\theta_{0}-\omega$.

As a conclusion, we have that $\varphi\left(\theta_{0}\right)=\eta\left(\theta_{0}\right)$ if and only if $b=b^{*}$, and that the curves are disjoint if $b<b^{*}$. If $b>b^{*}$ the two curves meet transversally in two points. Finally, $\eta$ is selfsymmetric by the symmetry of the map $S:(x, \theta) \mapsto(-x, \theta+\pi)$. This symmetry also implies that if we replace $\varphi_{0}(\theta)$ by $\gamma_{0}(\theta)=-\frac{\pi}{2}+b \sin (\theta-\omega)$ the distance between $\gamma_{0}$ and $\eta$ is the same as the distance between $\varphi_{0}$ and $\eta$.
Lemma 2. We recall that $\varphi_{0}(\theta)=\pi / 2+b \sin (\theta-\omega)$ and that $\gamma_{0}(\theta)=-\frac{\pi}{2}+b \sin (\theta-\omega)$. Let us define $\lambda_{0}=\varphi_{0}-\eta$, and

$$
\begin{aligned}
\varphi_{n+1}(\theta) & =h_{a}\left(\varphi_{n}(\theta-\omega)\right)+b \sin (\theta-\omega) \\
\gamma_{n+1}(\theta) & =h_{a}\left(\gamma_{n}(\theta-\omega)\right)+b \sin (\theta-\omega) \\
\lambda_{n+1}(\theta) & =h_{a}\left(\lambda_{n}(\theta-\omega)+\eta(\theta-\omega)\right)-a \eta(\theta-\omega) .
\end{aligned}
$$

for all $n \geq 0$. Then, if $a>0$, we have that

1. $\lambda_{n}=\varphi_{n}-\eta$ for any $n \geq 0$,
2. if $0 \leq b \leq b^{*}$, then $\lambda_{n} \geq 0$,
3. the sequences $\left\{\varphi_{n}\right\}_{n \geq 0}$ and $\left\{\lambda_{n}\right\}_{n \geq 0}$ are decreasing and $\left\{\gamma_{n}\right\}_{n \geq 0}$ is increasing.

## Proof:

1. By definition it is true for $n=0$, so we will proceed by induction: we assume that it is true for $n-1$, and we show it holds for $n$.

$$
\begin{aligned}
\lambda_{n}(\theta) & =h_{a}\left(\lambda_{n-1}(\theta-\omega)+\eta(\theta-\omega)\right)-a \eta(\theta-\omega) \\
& =h_{a}\left(\varphi_{n-1}(\theta-\omega)\right)-\eta(\theta)+b \sin (\theta-\omega)=\varphi_{n}(\theta)-\eta(\theta)
\end{aligned}
$$

where we have used that $\eta(\theta)=a \eta(\theta-\omega)+b \sin (\theta-\omega)$.
2. This follows from Lemma 1.
3. Let us start by showing that $\varphi_{1} \leq \varphi_{0}$ :

$$
\varphi_{1}(\theta)=h_{a}\left(\varphi_{0}(\theta-\omega)\right)+b \sin (\theta-\omega) \leq \frac{\pi}{2}+b \sin (\theta-\omega)=\varphi_{0}(\theta)
$$

Now, using that $h_{a}$ is an increasing function we can easily check that if $\varphi_{n-1} \leq \varphi_{n-2}$ then $\varphi_{n} \leq \varphi_{n-1}$ and the result on $\left\{\varphi_{n}\right\}_{n}$ follows by induction. To show that $\left\{\gamma_{n}\right\}_{n}$ is increasing we can use a similar proof, or simply to recall that $\gamma_{n}$ is the mirror image of $\varphi_{n}$ by the symmetry $S$. Finally, as $\lambda_{n}=\varphi_{n}-\eta$ we have that $\left\{\lambda_{n}\right\}_{n}$ is decreasing.
For each $b \in\left[0, b^{*}\right]$ we consider the following subsets of $\mathbb{R} \times \mathbb{T}^{1}$,

$$
\begin{aligned}
& A_{+}=\left\{(x, \theta) \in \mathbb{R} \times \mathbb{T}^{1} \mid \eta(\theta) \leq x \leq \pi / 2+b \sin (\theta-\omega)\right\} \\
& A_{-}=\left\{(x, \theta) \in \mathbb{R} \times \mathbb{T}^{1} \mid-\pi / 2+b \sin (\theta-\omega) \leq x \leq \eta(\theta)\right\}
\end{aligned}
$$

Note that $S\left(A_{+}\right)=A_{-}$and viceversa. In what follows we focus on the set $A_{+}$, and the results will be translated to $A_{-}$by the symmetry $S$.

Lemma 3. If $a>0$ and $|b| \leq b^{*}$, then $F_{a, b}\left(A_{+}\right) \subset A_{+}$.

Proof: If $(x, \theta) \in A_{+}$we can distinguish two cases:

1. If $x \geq \pi /(2 a)$ then $F_{a, b}(x, \theta)=(\pi / 2+b \sin (\theta), \theta+\omega) \in A_{+}$.
2. If $\eta(\theta) \leq x \leq \pi /(2 a)$ then $\eta(\theta+\omega)=a \eta(\theta)+b \sin (\theta) \leq a x+b \sin (\theta) \leq \pi / 2+b \sin (\theta)$, and this also implies that $F_{a, b}(x, \theta) \in A_{+}$.

Then, we have a compact invariant set given by

$$
\Lambda_{+}=\bigcap_{n=0}^{\infty} F_{a, b}^{n}\left(A_{+}\right)
$$

As, for each $\theta,\left\{\varphi_{n}(\theta)\right\}_{n}$ is a decreasing sequence that is bounded from below (by $\eta(\theta)$ ), it has a limit. Therefore, the pointwise limit of $\left\{\varphi_{n}\right\}_{n}$ is an at least upper semicontinuous invariant curve $\varphi_{\infty}$ contained in $\Lambda_{+}$. Moreover, $(x, \theta) \in \Lambda_{+}$if and only if $\eta(\theta) \leq x \leq \varphi_{\infty}(\theta)$.

As $\left\{\lambda_{n}(\theta)\right\}_{n}$ is decreasing and bounded from below by 0 , it has a limit $\lambda(\theta)=\lim _{n \rightarrow \infty} \lambda_{n}(\theta)$ that is at least upper semi-continuous. We have

$$
\lambda_{n+1}(\theta)=\left\{\begin{array}{ccc}
\pi / 2-a \eta(\theta-\omega) & \text { if } & \lambda_{n}(\theta-\omega)+\eta(\theta-\omega) \geq \pi /(2 a),  \tag{9}\\
a \lambda_{n}(\theta-\omega) & \text { if } & \lambda_{n}(\theta-\omega)+\eta(\theta-\omega) \leq \pi /(2 a) .
\end{array}\right.
$$

Now, we define

$$
I_{n}=\left\{\theta \in \mathbb{T}^{1} \mid \lambda_{n}(\theta-\omega)+\eta(\theta-\omega) \geq \pi /(2 a)\right\} .
$$

Note that if $\theta \in I_{n+1}$ then

$$
\lambda_{n}(\theta-\omega)+\eta(\theta-\omega) \geq \lambda_{n+1}(\theta-\omega)+\eta(\theta-\omega) \geq \pi /(2 a)
$$

which implies that $I_{n+1} \subset I_{n}$. Then $I=\bigcap_{i=0}^{\infty} I_{n} \neq \emptyset$ is a compact set such that if $\theta \in I$ then

$$
\lambda(\theta)=\frac{\pi}{2}-a \eta(\theta-\omega)
$$

### 2.1 A strange non-chaotic attractor

In this section we assume that $a>1$ is fixed and that $b=b^{*}$. As a side comment, it easily follows from (7) that $b^{*} \rightarrow \pi / 2$ when $a \rightarrow \infty$. The goal here is to show that the repelling continuous invariant curve $\eta$ and the two attracting invariant curves $\varphi_{\infty}(\theta)$ and $\gamma_{\infty}(\theta)$,

$$
-\frac{\pi}{2}+b \sin (\theta-\omega) \leq \gamma_{\infty}(\theta) \leq \eta(\theta) \leq \varphi_{\infty}(\theta) \leq \frac{\pi}{2}+b \sin (\theta-\omega)
$$

are such that $\gamma_{\infty}(\theta)$ is lower semicontinuous, $\varphi_{\infty}(\theta)$ is upper semicontinuous and all three curves are different. Each of the curves $\varphi_{\infty}, \gamma_{\infty}$ intersects $\eta$ on a dense set of zero measure of values of $\theta$. This means that the curves $\varphi_{\infty}, \gamma_{\infty}$ are not continuous everywhere.

Lemma 4. Let us define $\theta_{0}$ as the only zero of $\lambda_{0}=\varphi_{0}-\eta$. Then, $\lambda_{n}$ has exactly $n+1$ zeros at $\theta_{0}, \theta_{0}+\omega, \ldots, \theta_{0}+n \omega$.

Proof: Given any value $n \geq 1$, we note that

$$
\begin{aligned}
F_{a, b^{*}(a)}^{n-1}\left(\eta\left(\theta_{0}-n \omega\right), \theta_{0}-n \omega\right) & =\left(\eta\left(\theta_{0}-\omega\right), \theta_{0}-\omega\right), \\
F_{a, b^{*}(a)}^{n-1}\left(\varphi_{0}\left(\theta_{0}-n \omega\right), \theta_{0}-n \omega\right) & =\left(\varphi_{n-1}\left(\theta_{0}-\omega\right), \theta_{0}-\omega\right) .
\end{aligned}
$$

As $\eta\left(\theta_{0}-\omega\right)=\pi /(2 a)$ and $\varphi_{n-1}\left(\theta_{0}-\omega\right) \geq \eta\left(\theta_{0}-\omega\right)$ we have that

$$
F_{a, b^{*}(a)}\left(\eta\left(\theta_{0}-\omega\right), \theta_{0}-\omega\right)=F_{a, b^{*}(a)}\left(\varphi_{n-1}\left(\theta_{0}-\omega\right), \theta_{0}-\omega\right)
$$

and this implies that $\lambda_{n}\left(\theta_{0}\right)=\varphi_{n}\left(\theta_{0}\right)-\eta\left(\theta_{0}\right)$ is zero. To complete the existence of the $n+1$ zeros note that, if $\theta$ is a zero of $\lambda_{n}$ then from (9) it follows that $\lambda_{n+1}(\theta+\omega)=a \lambda_{n}(\theta)=0$.

To see that these are the only zeroes, we start from the fact that $\lambda_{0}$ has only one zero and we proceed by induction. If the result is true for $n \geq 0$ and we assume that $\lambda_{n+1}$ has an extra zero $\theta_{1}$ then from (9) there are two possibilities: i) $\pi / 2-a \eta\left(\theta_{1}-\omega\right)=0$, or ii) $\lambda_{n}\left(\theta_{1}-\omega\right)=0$. For the first option, note that $\pi / 2-a \eta\left(\theta_{1}-\omega\right)=\lambda_{0}\left(\theta_{1}\right)$ which contradicts the fact that $\lambda_{0}$ has
$\theta_{0}$ as its only zero. The second option implies that $\lambda_{n}$ has an extra zero. We can repeat this reasoning until we show (either using i) or ii)) that $\lambda_{0}$ has an extra zero.

Moreover, we have that $\lambda_{0}\left(\theta_{0}\right)=0$ and $\lambda_{0}(\theta)>0$, for all $\theta \in \mathbb{T}^{1} \backslash\left\{\theta_{0}\right\}$. Then, $\lambda_{n}(\theta)$ has $n+1$ double zeros at $\theta_{0}, \theta_{0}+\omega, \ldots, \theta_{0}+n \omega$. This implies that $\lambda(\theta)$ is equal to zero in all the forward orbit of $\theta_{0}$. Note that, when $\lambda_{n}(\theta)$ is very small it means that $\varphi_{n}(\theta)$ is very close to $\eta(\theta)$ so it implies that $\lambda_{n+1}(\theta+\omega)=a \lambda_{n}(\theta)$. This justifies to define, by continuity, that $\lambda_{n+1}(\theta+\omega) / \lambda_{n}(\theta)=a$ when $\lambda_{n}(\theta)=0$. Therefore, we define

$$
\psi_{n}(\theta)=\left\{\begin{array}{cc}
\frac{\lambda_{n+1}(\theta+\omega)}{\lambda_{n}(\theta)} & \text { if } \lambda_{n}(\theta) \neq 0  \tag{10}\\
a & \text { if } \lambda_{n}(\theta)=0
\end{array}\right.
$$

We note that $\psi_{n}(\theta)>0$ for all $\theta \in \mathbb{T}^{1} \backslash\left\{\theta_{0}-\omega\right\}$.
Lemma 5. For $a \geq 1$ and every $n \geq 1$,

$$
0 \leq \psi_{n-1}(\theta) \leq \psi_{n}(\theta) \leq a,
$$

for all $\theta \in \mathbb{T}^{1}$.
Proof: It is clear that

$$
\psi_{n}(\theta)=\left\{\begin{array}{cl}
\frac{\lambda_{0}(\theta+\omega)}{\lambda_{n}(\theta)} & \text { if } \lambda_{n}(\theta)+\eta(\theta) \geq \pi /(2 a)  \tag{11}\\
a & \text { if } \lambda_{n}(\theta)+\eta(\theta) \leq \pi /(2 a)
\end{array}\right.
$$

Now we have to distinguish several cases, depending on the pieces of the map $h_{a}$.

1) $\lambda_{n-1}(\theta)+\eta(\theta) \geq \pi /(2 a)$.
1.a) $\lambda_{n}(\theta)+\eta(\theta) \geq \pi /(2 a)$. In this case,

$$
\psi_{n-1}(\theta)=\frac{\lambda_{0}(\theta+\omega)}{\lambda_{n-1}(\theta)} \quad \text { and } \quad \psi_{n}(\theta)=\frac{\lambda_{0}(\theta+\omega)}{\lambda_{n}(\theta)}
$$

Then, as $\lambda_{n}(\theta) \leq \lambda_{n-1}(\theta)$ we have that $\psi_{n}(\theta) \geq \psi_{n-1}(\theta)$.
1.b) $\lambda_{n}(\theta)+\eta(\theta) \leq \pi /(2 a)$. Here,

$$
\psi_{n-1}(\theta)=\frac{\lambda_{0}(\theta+\omega)}{\lambda_{n-1}(\theta)} \quad \text { and } \quad \psi_{n}(\theta)=a
$$

Now, because $\lambda_{n-1}(\theta)+\eta(\theta) \geq \pi /(2 a)$ we obtain

$$
\lambda_{n-1}(\theta) \geq \frac{\pi}{2 a}-\eta(\theta)=\frac{1}{a}\left(\frac{\pi}{2}-a \eta(\theta)\right)=\frac{1}{a} \lambda_{0}(\theta+\omega),
$$

and, therefore, $\psi_{n}(\theta) \geq \psi_{n-1}(\theta)$.
2) $\lambda_{n-1}(\theta)+\eta(\theta) \leq \pi /(2 a)$. In this situation, as we also have $\lambda_{n}(\theta)+\eta(\theta) \leq \pi /(2 a)$, we conclude that $\psi_{n}(\theta)=\psi_{n-1}(\theta)=a$.

To complete the proof, we use (11):

$$
\psi_{n}(\theta)=\frac{\lambda_{0}(\theta+\omega)}{\lambda_{n}(\theta)} \leq \frac{\lambda_{0}(\theta+\omega)}{\frac{\pi}{2 a}-\eta(\theta)}=\frac{\lambda_{0}(\theta+\omega)}{\lambda_{0}(\theta+\omega) / a}=a .
$$

Lemma 6. The sequence $\left\{\psi_{n}\right\}_{n}$ converges pointwise to a measurable and integrable function $\psi$, and

$$
\int_{0}^{2 \pi} \log \psi(\theta) d \theta \leq 0
$$

Proof: Let us see first that the functions $\log \psi_{n}$ are integrable. As Lemma 5 shows that $0 \leq \psi_{n} \leq a$ and $\psi_{0} \leq \psi_{n}$, we only have to see that $\log \psi_{0}$ is integrable. This follows from the fact that

$$
\psi_{0}(\theta)=\left\{\begin{array}{cl}
\frac{\pi / 2-a \eta(\theta)}{\pi / 2-a \eta(\theta-\omega)} & \text { if } \lambda_{0}(\theta)+\eta(\theta) \geq \pi /(2 a), \\
a & \text { if } \lambda_{0}(\theta)+\eta(\theta) \leq \pi /(2 a),
\end{array}\right.
$$

that is, $\psi_{0}$ is well defined, continuous and it only takes the value 0 for a single value of $\theta=\theta_{0}$, which is a zero of multiplicity two. This implies that $\log \psi_{0}$ is integrable and, therefore, all $\log \psi_{n}$ are integrable. Moreover, as $\left\{\lambda_{n}(\theta)\right\}_{n}$ is a decreasing sequence for all $\theta$ we have that, for a set of values of $\theta$ of full measure (those for which $\lambda_{n}(\theta) \neq 0$ ),

$$
\psi_{n}(\theta) \leq \frac{\lambda_{n}(\theta+\omega)}{\lambda_{n}(\theta)}
$$

Hence, for those values of $\theta$,

$$
\log \psi_{n}(\theta) \leq \log \lambda_{n}(\theta+\omega)-\log \lambda_{n}(\theta)
$$

As $\log \lambda_{n}$ is an integrable function (the zeroes of $\lambda_{n}$ are of multiplicity 2 ), we have that

$$
\int_{0}^{2 \pi} \log \psi_{n}(\theta) d \theta \leq \int_{0}^{2 \pi} \log \lambda_{n}(\theta+\omega) d \theta-\int_{0}^{2 \pi} \log \lambda_{n}(\theta) d \theta=0
$$

Finally, using the dominated convergence theorem, the statement follows.
Theorem 2.1. For each $\theta \in \mathbb{T}^{1}$, let us consider the sequence $\left\{\varphi_{n}(\theta)\right\}_{n}$.
a) The sequence $\left\{\varphi_{n}(\theta)\right\}_{n}$ has a limit for each $\theta$, and the function $\varphi_{\infty}$ defined as the point limit

$$
\varphi_{\infty}(\theta)=\lim _{n \rightarrow \infty} \varphi_{n}(\theta)
$$

is upper semicontinuous and invariant.
b) The set $A=\left\{\theta \in \mathbb{T}^{1}\right.$ such that $\left.\varphi_{\infty}(\theta)=\eta(\theta)\right\}$ is dense and it has zero Lebesgue measure.

## Proof:

a) We note that $\varphi_{\infty}$ is well-defined since, for each $\theta$, the sequence $\left\{\varphi_{n}(\theta)\right\}_{n}$ is decreasing and bounded from below. This also implies that the function $\varphi_{\infty}$ is upper semicontinuous. The invariance follows from the definition of the sequence $\varphi_{n}$,

$$
\varphi_{n+1}(\theta)=h_{a}\left(\varphi_{n}(\theta-\omega)\right)+b^{*} \sin (\theta-\omega)
$$

b) As $\theta_{0}$ belongs to $A$, all the points $\theta_{0}+k \omega, k \in \mathbb{N}$, belong to $A$. Since $\omega /(2 \pi)$ is irrational, $A$ is dense in $\mathbb{T}^{1}$. To show that $A$ has zero Lebesgue measure we will show first that

$$
A \backslash\left\{\theta_{0}-\omega\right\} \subset\left\{\theta \in \mathbb{T}^{1} \text { such that } \psi(\theta)=a\right\}
$$

To see it, let us select an arbitrary point $\theta \in A, \theta \neq \theta_{0}-\omega$. Then, $\varphi_{\infty}(\theta)=\eta(\theta)<\pi /(2 a)$. Therefore, there exists a value $n_{0}$ (depending on $\theta$ ) such that $\varphi_{n}(\theta)<\pi /(2 a)$ for $n \geq n_{0}$. This implies that $\lambda_{n+1}(\theta+\omega)=a \lambda_{n}(\theta)$ and this is equivalent to $\psi_{n}(\theta)=a$ for all $n \geq n_{0}$ and, hence, $\psi(\theta)=a$. To complete this part, let us assume that the Lebesgue measure of $A$ is not zero. Then, by ergodicity, $A$ has to have total measure. This implies that the measure of the set of values of $\theta$ such that $\psi(\theta)=a$ has total measure and then, since $a>1$,

$$
\int_{0}^{2 \pi} \log \psi(\theta) d \theta=\log a>0
$$

which contradicts Lemma 6 . Therefore, $A$ has zero measure.
Let us now focus on Lyapunov exponents. It is well known that the Lyapunov exponent of an invariant curve $\phi$ of a smooth map is

$$
\begin{equation*}
\Lambda=\frac{1}{2 \pi} \int_{0}^{2 \pi} \log \left(\left|\partial_{x} f(\phi(\theta), \theta)\right|\right) d \theta \tag{12}
\end{equation*}
$$

In our case, $f(x, \theta)=h_{a}(x)+b \sin \theta$ and $h_{a}$ is piecewise linear, and its derivative is well defined except in two points, $x= \pm \pi /(2 a)$. Moreover, the attracting curve $\varphi_{\infty}$ crosses $x=\pi /(2 a)$, which means that we cannot directly write the formula (12) since the derivative w.r.t. $x$ is not well defined at $x=\pi /(2 a)$. On the other hand, $h_{a}$ has left and right derivatives at $x=\pi /(2 a)$ equal to $a$ and 0 respectively. The value $a$ corresponds to an expansion $(a>1)$ while 0 corresponds to a compression. For the moment being, let us define the derivative of $h_{a}$ at $x=\pi /(2 a)$ as $a$ and let us compute the Lyapunov exponent of $\varphi_{\infty}$. We note that, with this assumption, the curve $\varphi_{\infty}$ will be considered "as repelling as possible".

Proposition 1. The Lyapunov exponent of $\varphi_{\infty}$ is $-\infty$.

Proof: Let us first show that the measure of the set

$$
B=\left\{\theta \in \mathbb{T}^{1} \text { such that } \varphi_{\infty}(\theta)>\pi /(2 a)\right\}
$$

is strictly positive. To see it, let us assume that it is zero. Then, we consider $\lambda=\varphi_{\infty}-\eta$ and let us show that all the Fourier coefficients of $\lambda, \lambda^{(k)}$, are zero:

$$
\lambda^{(k)}=\frac{1}{2 \pi} \int_{0}^{2 \pi} \lambda(\theta) e^{-k \omega \mathbf{i}} d \theta=\frac{1}{2 \pi} \int_{\mathbb{T}^{1} \backslash B} \lambda(\theta) e^{-k \omega \mathbf{i}} d \theta
$$

As $\varphi_{\infty}(\theta) \leq \pi /(2 a)$ for $\theta \in \mathbb{T}^{1} \backslash B$, the functions $\varphi_{\infty}$ and $\eta$ take values on $[-\pi /(2 a), \pi /(2 a)]$ and, as they are invariant, we have that $\lambda(\theta+\omega)=a \lambda(\theta)$, which forces that $\lambda^{(k)} e^{k \omega \mathbf{i}}=a \lambda^{(k)}$ and, as $a>1$, we conclude $\lambda^{(k)}=0$ and then $\lambda=0$ a.e., which contradicts $b$ ).

Now,

$$
\begin{aligned}
\Lambda & =\frac{1}{2 \pi} \int_{0}^{2 \pi} \log \left(h_{a}^{\prime}\left(\varphi_{\infty}(\theta)\right)\right) d \theta \\
& =\frac{1}{2 \pi}\left(\int_{B} \log \left(h_{a}^{\prime}\left(\varphi_{\infty}(\theta)\right)\right) d \theta+\int_{\mathbb{T}^{1} \backslash B} \log \left(h_{a}^{\prime}\left(\varphi_{\infty}(\theta)\right)\right) d \theta\right) \\
& =\frac{1}{2 \pi}\left((-\infty) \operatorname{meas}(B)+(\log a) \operatorname{meas}\left(\mathbb{T}^{1} \backslash B\right)\right)=-\infty,
\end{aligned}
$$

because meas $(B)>0$.
Remark 1. Note that the value assigned to $h^{\prime}(\pi /(2 a))$ is irrelevant as long as it is finite.
Proposition 2. There exists a set $E \subset \mathbb{T}^{1}$ with total Lebesgue measure such that, for any point $\left(x_{0}, \theta_{0}\right) \in \mathbb{R} \times E$ such that $x_{0}>\eta\left(\theta_{0}\right)$, the orbit $\left(x_{n+1}, \theta_{n+1}\right)=F_{a, b^{*}}\left(x_{n}, \theta_{n}\right)$ satisfies that there exists a natural value $m=m\left(x_{0}, \theta_{0}\right)$ such that $x_{m}=\varphi_{\infty}\left(\theta_{m}\right)$. That is, the orbit starting at $\left(x_{0}, \theta_{0}\right)$ arrives to the attracting curve $\varphi_{\infty}$ in a finite number of iterates.

Proof: Let us choose an arbitrary point $\left(x_{0}, \theta_{0}\right)$ (for the moment being, $\theta_{0} \in \mathbb{T}^{1}$ ). As $x_{0}>\eta\left(\theta_{0}\right)$, we have that $\eta\left(\theta_{n}\right)<x_{n}(n \geq 0)$. Let us define the rotation $T_{\omega}: \mathbb{T}^{1} \rightarrow \mathbb{T}^{1}$ as $T_{\omega}(\theta)=\theta+\omega$. We recall that we have defined the set $B$ as

$$
B=\left\{\theta \in \mathbb{T}^{1} \text { such that } \varphi_{\infty}(\theta)>\pi /(2 a)\right\}
$$

Then, we define the set $C$ as

$$
C=\bigcup_{n \in \mathbb{N}} T_{\omega}^{-n}(B) .
$$

As the Lebesgue measure is invariant and ergodic by $T_{\omega}$ (recall that $\omega$ is irrational), using Theorem 1.5 in [Wal82] we obtain that $C$ has total measure. Now, we distinguish several cases.
a) Case $\varphi_{\infty}\left(\theta_{0}\right) \leq x_{0}, \theta_{0} \in C$.

As $\theta_{0} \in C$, there exists a value $n_{0} \in \mathbb{N}$ such that $\theta_{n_{0}}=T_{\omega}^{n_{0}}\left(\theta_{0}\right) \in B$. Then, we have that $\pi /(2 a)<\varphi_{\infty}\left(\theta_{n_{0}}\right) \leq x_{n_{0}}$, and this implies that $\varphi_{\infty}\left(\theta_{n_{0}+1}\right)=x_{n_{0}+1}$.
b) Case $\eta\left(\theta_{0}\right)<x_{0} \leq \varphi_{\infty}\left(\theta_{0}\right), \theta_{0} \in D=\left\{\theta \in \mathbb{T}^{1}\right.$ such that $\left.\varphi_{\infty}(\theta) \neq \eta(\theta)\right\}$.

We note that this is a total measure set. First, let us note that there exists $n_{0} \in \mathbb{N} \cup\{0\}$ such that $x_{n_{0}}>\pi /(2 a)$ (otherwise, we would have that $x_{n}-\eta\left(\theta_{n}\right)=a^{n}\left(x_{0}-\eta\left(\theta_{0}\right)\right)$ and then $x_{n}$ would become unbounded). Then, $x_{n_{0}+1}=\varphi_{0}\left(\theta_{n_{0}+1}\right) \leq \varphi_{\infty}\left(\theta_{n_{0}+1}\right)$ and, on the other hand, $\varphi_{0}\left(\theta_{n_{0}+1}\right) \geq \varphi_{\infty}\left(\theta_{n_{0}+1}\right)$, which implies $x_{n_{0}+1}=\varphi_{\infty}\left(\theta_{n_{0}+1}\right)$.

To complete the proof we define $E=C \cap D$, which is also of full measure.
Remark 2. As we have mentioned before, we can use the symmetry $S$ to transport these results to the set $A_{-}$. So, we can define $\varphi^{-}(\theta)=-\varphi_{\infty}(\theta+\pi)$ and then it is a lower semicontinuous function such that the set $\left\{\theta \in \mathbb{T}^{1}\right.$ such that $\left.\gamma_{\infty}(\theta)=\eta(\theta)\right\}$ is dense and it has zero Lebesgue measure. Moreover, the Lyapunov exponent of $\gamma_{\infty}$ is $-\infty$ and the orbit starting at an arbitrary point $\left(x_{0}, \theta_{0}\right)$ such that $x_{0}<\eta\left(\theta_{0}\right)$ arrives to $\gamma_{\infty}$ in a finite number of iterates.

### 2.2 Non-smooth pitchfork bifurcation

In this section we still assume that $a>1$. The goal here is to show that, for $0<b<b^{*}$, the map $F_{a, b}$ has three continuous invariant curves, one repelling and two attracting, and for $b>b^{*}$, the map $F_{a, b}$ has only one invariant curve which is continuous and attracting. We have already seen in Section 2.1 that, when $b=b^{*}$, the two attracting curves that existed for $0<b<b^{*}$ collide with the repelling curve creating a strange non-chaotic attractor.

### 2.2.1 Three invariant curves

Let us start with the case $b<b^{*}$. The existence of a repelling curve $\eta$ has already been shown at the beginning of Section 2. Note that we only need to show the existence of an attracting invariant curve above $\eta$, since by means of the symmetry $S$ we will obtain immediately the existence of another attracting invariant curve below $\eta$.

Proposition 3. If $0<b<b^{*}$ there exists a unique continuous attracting invariant curve $\varphi_{\infty}$ such that $\eta(\theta)<\varphi_{\infty}(\theta) \leq \varphi_{0}(\theta)=\pi / 2+b \sin (\theta-\omega)$ for all $\theta \in \mathbb{T}^{1}$.

Proof: Let us define $\delta_{b}>0$ as (see Lemma 1)

$$
\delta_{b}=\min _{\theta \in \mathbb{T}}\left(\frac{\pi}{2 a}-\eta\right)=\frac{\pi}{2 a}-\frac{b}{\sqrt{1+a^{2}-2 a \cos \omega}}>0
$$

and let us also define $\eta_{0}(\theta)=\eta(\theta)+\delta_{b}$. Obviously, $\eta_{0}(\theta) \in[-\pi /(2 a), \pi /(2 a)]$. Moreover, Lemma 1 also implies that $\eta_{0}(\theta) \leq \varphi_{0}(\theta)$. Next, let us define the sequences

$$
\eta_{n}(\theta)=h_{a}\left(\eta_{n-1}(\theta-\omega)\right)+b \sin (\theta-\omega), \quad \varphi_{n}(\theta)=h_{a}\left(\varphi_{n-1}(\theta-\omega)\right)+b \sin (\theta-\omega),
$$

for $n \geq 1$ and for all $\theta \in \mathbb{T}$. Let us see that $\eta_{0}(\theta) \leq \eta_{1}(\theta)$ :

$$
\begin{aligned}
\eta_{1}(\theta) & =h_{a}\left(\eta_{0}(\theta-\omega)\right)+b \sin (\theta-\omega)=a \eta_{0}(\theta-\omega)+b \sin (\theta-\omega) \\
& =a \eta(\theta-\omega)+a \delta_{b}+b \sin (\theta-\omega)=\eta(\theta)+\delta_{b}+(a-1) \delta_{b}>\eta_{0}(\theta) .
\end{aligned}
$$

The monotonicity of $h_{a}$ implies that $\eta_{n}(\theta) \geq \eta_{n-1}(\theta)$ and $\eta_{n}(\theta) \leq \varphi_{n}(\theta)$ for all $n \geq 1$ and $\theta \in \mathbb{T}$. Therefore, the sequence $\left\{\eta_{n}\right\}_{n \geq 0}$ is increasing and bounded from above, which implies that it is pointwise convergent to a lower semicontinuous invariant curve $\eta_{\infty}$. We have seen in Lemma 2 that the sequence $\left\{\varphi_{n}\right\}_{n \geq 0}$ is decreasing and bounded from below, which implies that it is pointwise convergent to an upper semicontinuous invariant curve $\varphi_{\infty}$.

To complete the proof, we will show that $\eta_{\infty}=\varphi_{\infty}$. To this end, let us choose a fixed $\theta_{0} \in \mathbb{T}$ and let us consider the images of $\left(\eta_{0}\left(\theta_{0}\right), \theta_{0}\right)$ by the map $F_{a, b}$. It is clear that $\eta_{0}\left(\theta_{0}\right) \in$ $[-\pi /(2 a), \pi /(2 a)]$, and the iterates of this point are defined as

$$
F_{a, b}^{n}\left(\eta_{0}\left(\theta_{0}\right), \theta_{0}\right)=\left(\eta_{n}\left(\theta_{0}+n \omega\right), \theta_{0}+n \omega\right) .
$$

Let us see that, after a suitable number $\left(n_{0}\right)$ of iterates of the map $F_{a, b}$, the first component satisfies $\eta_{n_{0}}\left(\theta_{0}+n_{0} \omega\right)>\pi /(2 a)$. So, assume that this is not true. This would imply that, for all $n$, we have $\eta_{n}\left(\theta_{0}+n \omega\right)=\eta(\theta+n \omega)+a^{n} \delta_{b}$ which is absurd since $a^{n} \delta_{b}$ is not bounded. Therefore, by continuity of the maps $\eta_{n}$, there exists a small open neighbourhood $I_{0}$ of $\theta_{0}$ such that, for
all $\theta \in I_{0}$, we have that $\eta_{n_{0}}\left(\theta+n_{0} \omega\right)>\pi /(2 a)$. The definition of the map $F_{a, b}$ implies that the next image of these points satisfies that $\eta_{n_{0}+1}\left(\theta+\left(n_{0}+1\right) \omega\right)=\varphi_{0}\left(\theta+\left(n_{0}+1\right) \omega\right)$ for all $\theta \in I_{0}$ which means that, on the set $I_{0}$, the two limit curves $\eta_{\infty}$ and $\varphi_{\infty}$ coincide. Finally, as the two curves are invariant and coincide on an open set, they must be equal. Moreover, this attracting invariant curve is continuous.

### 2.2.2 One invariant curve

Here we consider the case $b>b^{*}$. We proceed in a similar way as in the previous section.
Proposition 4. If $a \geq 1$ and $b>b^{*}$ there exists a unique, self-symmetric continuous attracting invariant curve.

Proof: Let us consider the sequences $\left\{\varphi_{n}\right\}_{n \geq 0}$ and $\left\{\gamma_{n}\right\}_{n \geq 0}$ that have been defined in Lemma 2, with pointwise limits $\varphi_{\infty}$ and $\gamma_{\infty}$ respectively. We recall that $\varphi_{\infty}$ and $\gamma_{\infty}$ are related by the symmetry $S$. To show that $\varphi_{\infty}=\gamma_{\infty}$ (this also implies that they are continuous and selfsymmetric) we will show that $\varphi_{n}$ coincides with $\gamma_{n}$ for $n$ large enough. To this end, we will prove first that there exists a point $\theta_{0}$ such that, for a suitable $n_{1}, \varphi_{n_{1}}$ coincides with $\gamma_{n_{1}}$ on an open neighbourhood of $\theta_{0}$. Then, we will show that this implies that $\varphi_{n_{2}}=\gamma_{n_{2}}$ for a suitable $n_{2}>n_{1}$. Let us see the details.

Let $\theta_{0}$ be the value that minimizes $\varphi_{0}(\theta)-\eta(\theta)$. According to Lemma 1,

$$
\begin{equation*}
\varphi_{0}\left(\theta_{0}\right)-\eta\left(\theta_{0}\right)=\frac{\pi}{2}-\frac{a b}{\sqrt{1+a^{2}-2 a \cos \omega}}=a \min _{\theta \in \mathbb{T}}\left(\frac{\pi}{2 a}-\eta(\theta)\right), \tag{13}
\end{equation*}
$$

and, if $b>b^{*}$, we have that $\varphi_{0}\left(\theta_{0}\right)-\eta\left(\theta_{0}\right)<0$. Let us see that there exists $n_{0} \geq 0$ such that $\varphi_{n_{0}}\left(\theta_{0}+n_{0} \omega\right)<-\pi /(2 a)$. First, from (13) we see that

$$
\varphi_{0}\left(\theta_{0}\right)-\eta\left(\theta_{0}\right) \leq a\left(\frac{\pi}{2 a}-\eta\left(\theta_{0}\right)\right)<\frac{\pi}{2 a}-\eta\left(\theta_{0}\right)
$$

which implies that $\varphi_{0}\left(\theta_{0}\right)<\pi /(2 a)$. If $\varphi_{0}\left(\theta_{0}\right)<-\pi /(2 a)$ then $n_{0}=0$. Otherwise, let us consider the points $\varphi_{n}\left(\theta_{0}+n \omega\right)$ and note that they cannot be inside $[-\pi /(2 a), \pi /(2 a)]$ for all $n$ because, if they are inside this interval they must satisfy that $\varphi_{n}\left(\theta_{0}+n \omega\right)-\eta\left(\theta_{0}+n \omega\right)=a^{n}\left(\varphi_{0}\left(\theta_{0}\right)-\eta\left(\theta_{0}\right)\right)$ which means that the distance from $\varphi_{n}$ to $\eta$ is unbounded and this is impossible. Therefore, let us define $n_{0}$ as the first value of $n$ for which $\varphi_{n}\left(\theta_{0}+n \omega\right)$ is outside $[-\pi /(2 a), \pi /(2 a)]$. Hence, using that $\min (\pi /(2 a)-\eta(\theta))<0$ we have
$\varphi_{n_{0}}\left(\theta_{0}+n_{0} \omega\right)-\eta\left(\theta_{0}+n_{0} \omega\right)=a^{n_{0}+1} \min _{\theta \in \mathbb{T}}\left(\frac{\pi}{2 a}-\eta(\theta)\right) \leq \min _{\theta \in \mathbb{T}}\left(\frac{\pi}{2 a}-\eta(\theta)\right) \leq \frac{\pi}{2 a}-\eta\left(\theta_{0}+n_{0} \omega\right)$,
which implies that $\varphi_{n_{0}}\left(\theta_{0}+n_{0} \omega\right) \leq \pi /(2 a)$ and, as $\varphi_{n_{0}}\left(\theta_{0}+n_{0} \omega\right)$ is outside $[-\pi /(2 a), \pi /(2 a)]$, we conclude that $\varphi_{n_{0}}\left(\theta_{0}+n_{0} \omega\right)<-\pi /(2 a)$.

By continuity, there exists an open neighbourhood $I_{0}$ of $\theta_{0}$ such that $\varphi_{n_{0}}\left(\theta+n_{0} \omega\right)<-\pi /(2 a)$ for all $\theta \in I_{0}$. This means that $\varphi_{n_{0}+1}\left(\theta+\left(n_{0}+1\right) \omega\right)=\gamma_{0}\left(\theta+\left(n_{0}+1\right) \omega\right)$ for all $\theta \in I_{0}$. Defining $n_{1}=n_{0}+1$, this forces that $\varphi_{n_{1}}\left(\theta+n_{1} \omega\right)=\gamma_{0}\left(\theta+n_{1} \omega\right)=\varphi_{\infty}\left(\theta+n_{1} \omega\right)=\gamma_{\infty}\left(\theta+n_{1} \omega\right)$ for all $\theta \in I_{0}$. This implies, as the sequence $\left\{\gamma_{n}\right\}_{n}$ is increasing and bounded by $\gamma_{\infty}$, that $\varphi_{n_{1}}\left(\theta+n_{1} \omega\right)=\gamma_{n_{1}}\left(\theta+n_{1} \omega\right)$. Moreover, $\varphi_{n}(\theta+n \omega)=\gamma_{n}(\theta+n \omega)$ for $n \geq n_{1}$ and, as $\omega \notin 2 \pi \mathbb{Q}$, we conclude that there exists $n_{2} \geq n_{1}$ such that $\varphi_{n_{2}}=\gamma_{n_{2}}=\varphi_{\infty}=\gamma_{\infty}$.

### 2.3 The case $a \leq 1$

Let us focus first with the case $a<1, b \geq 0$. If $C(\mathbb{T}, \mathbb{R})$ denotes the vector space of continuous functions from $\mathbb{T}$ to $\mathbb{R}$ endowed with the sup norm, it is clear that an invariant curve of $F_{a, b}$ is a fixed point of the map $\mathcal{F}: C(\mathbb{T}, \mathbb{R}) \rightarrow C(\mathbb{T}, \mathbb{R})$ defined as

$$
\mathcal{F}(\varphi)(\theta)=h_{a}(\varphi(\theta-\omega))-b \sin (\theta-\omega) .
$$

It is immediate to check that $\mathcal{F}$ is a Lipschitz function with Lipschitz constant $a<1$ so it is contractive. Then, the Banach fixed point theorem implies that there exists a unique invariant curve in $C(\mathbb{T}, \mathbb{R})$.

Now let us consider the case $a=1$. This is a degenerate case, for instance for $0<b<b^{*}$ there is a continuous family of invariant curves: Lemma 1 implies that the invariant curve $\eta$ defined in (6) is at a finite distance $d_{b}$ of $x= \pm \pi / 2$ and then all the curves $\eta+\delta$ with $\delta$ a constant such that $|\delta| \leq d_{b}$ are also invariant. If $b=b^{*}$ then $\delta=0$ and this set of invariant curves reduces to a single one, $\eta$. The distance $d_{b}$ is attained at a some angle $\theta_{0}$. As $\eta+d_{b}$ is an invariant curve, we have that $\eta+d_{b} \leq \varphi_{\infty} \leq \varphi_{0}$. Moreover, $\varphi_{0}\left(\theta_{0}\right)=\eta\left(\theta_{0}\right)+d_{b}$ and then $\varphi_{\infty}\left(\theta_{0}\right)=\eta\left(\theta_{0}\right)+d_{b}$. To see that these are the only invariant curves let us discuss the different options separately.
$a=1,0<b<b^{*}$. Let us consider the sequence $\left\{\lambda_{n}\right\}_{n}$ defined in Lemma 2 and then the sequence $\left\{\psi_{n}\right\}_{n}$ defined in (10). Lemma 5 shows that the sequence $\left\{\psi_{n}\right\}_{n}$ is increasing and bounded by 1 . Moreover, as $\lambda_{n}(\theta)$ is never zero, (10) implies that

$$
\begin{equation*}
\psi_{n}(\theta)=\frac{\lambda_{n+1}(\theta+\omega)}{\lambda_{n}(\theta)} \leq 1, \tag{14}
\end{equation*}
$$

and, as $\lambda(\theta)=\lim _{n \rightarrow \infty} \lambda_{n}(\theta)$, then

$$
\begin{equation*}
\lambda(\theta+\omega) \leq \lambda(\theta) \quad \text { for all } \quad \theta \in \mathbb{T}^{1} . \tag{15}
\end{equation*}
$$

As $\lambda$ is an upper semicontinuous function, let us see that the previous inequality implies that $\lambda$ is a constant function. To this end, we choose two arbitrary values $\hat{\theta}$ and $\tilde{\theta}$ to show that $\lambda(\hat{\theta}) \leq \lambda(\tilde{\theta})$ : as $\omega \notin 2 \pi \mathbb{Q}$ we select an integer sequence $\left\{k_{n}\right\}_{n}, k_{n}>0$, such that $\left\{\hat{\theta}-k_{n} \omega\right\}_{n}$ tends to $\tilde{\theta}$. Using (15) we have that the sequence $\left\{\lambda\left(\hat{\theta}-k_{n} \omega\right)\right\}_{n}$ is increasing and $\lambda(\hat{\theta}) \leq \lambda\left(\hat{\theta}-k_{n} \omega\right)$ for all $n$. Therefore,

$$
\lambda(\hat{\theta}) \leq \lim _{n \rightarrow \infty} \lambda\left(\hat{\theta}-k_{n} \omega\right) \leq \lambda(\tilde{\theta}) .
$$

This scheme can also be used to show that $\lambda(\tilde{\theta}) \leq \lambda(\hat{\theta})$, to conclude that $\lambda$ is a constant function. Therefore, as we have seen above that $\lambda\left(\theta_{0}\right)=d_{b}$, we conclude that $\lambda \equiv d_{b}$. To finish this case, we note that any invariant curve must be either above or below $\eta$. If it is above, then it must be contained between $\varphi_{n}$ and $\eta$ for all $n$, so it must be between $\varphi_{\infty}$ and $\eta$. Therefore, it must be one of the invariant curves $\eta+\delta$ for $0 \leq \delta \leq d_{b}$. Finally, if the invariant curve is below $\eta$ we arrive to the same conclusion by using the symmetry $S$. $a=1, b=b^{*}$. The idea is to prove (15) and then the proof is exactly the same as in the previous item. The main difference to define $\psi_{n}$ is that we have to take into account that $\lambda$ can be zero. If $\theta$ is a value such that $\lambda(\theta)=0$, then $\lambda(\theta+\omega)=0$ and (15) holds. If $\lambda(\theta) \neq 0$ then there exists a $n_{0}$ such that if $n \geq n_{0}, \lambda_{n}(\theta)>0$. Therefore, (14) holds and then (15) holds.
$a=1, b>b^{*}$. This case is included in Proposition 4.

### 2.4 A codimension two bifurcation

It is interesting to compare these results with the ones in [JMAT18] for the model (3). To simplify this comparison, we focus on the case where $\omega$ is the golden mean, $\omega=\pi(\sqrt{5}-1)$. In the model (3), if the parameter $a>1$ is close enough to 1 , there is a smooth pitchfork bifurcation of invariant curves, going from three curves (two attracting and one repelling) for $b$ small that merge into one neutral curve when $b$ reaches a critical value $b^{*}$ that becomes attracting when $b>b^{*}$. This picture changes drastically when $a$ is large enough: there is still three smooth curves for $b$ small and one curve for $b$ large, but the neutral bifurcating curve does not look like a continuous curve but a SNA. The paper [JMAT18] contains a numerical computation of the pitchfork bifurcating invariant curve starting from low values of $a$ (where still is a smooth curve) and going up in $a$ from there. It turns out that these bifurcating curves becomes more and more wrinkled until the number of Fourier modes needed to approximate them becomes very large (near $5 \times 10^{7}$ ) and the computation is stopped. The last computed curve corresponds to $(a, b)=(5.348847,1.905990)$. The numerical simulations seem to show that, shortly after this point (i.e., for larger values of $a$ ), there exists critical values for the parameters, namely ( $\hat{a}, \hat{b}$ ) such that, from this point on, the smooth bifurcating curve is replaced by a SNA.

This situation is very similar to what happens in the model (1), where all the previous calculations are done explicitly in the previous sections. We have seen that

1. the critical value $(\hat{a}, \hat{b})$ is here $(1, \pi \sqrt{2-2 \cos \omega} / 2)$,
2. there exists a curve $\left(a, b^{*}(a)\right)$, defined for $a \geq 1$, such that, for $a>1$, the model has a SNA for each point of this curve,
3. for $a>1$, if we cross transversally the previous curve there is a nonsmooth pitchfork bifurcation of invariant curves,
4. for $0<a<1$ and $b<\hat{b}$, there is only a smooth attracting invariant curve that, if we increase the value of $a$ (and keeping the value of $b$ ), this curve becomes three invariant curves after crossing $a=1$.

These properties are also shown in Figure 2. We note that the critical value ( $\hat{a}, \hat{b}$ ) seems to play the same role in both models. It is remarkable that to unfold all the possible behaviours of the system around this point two parameters are needed.

## 3 A piecewise constant quasiperiodically forced map

Here we focus on the pointwise limit case when $a \rightarrow \infty$ (either in (2) or (3)), that is, the map $(\bar{x}, \bar{\theta})=F_{b}(x, \theta)$ defined as

$$
\left\{\begin{array}{l}
\bar{x}=h(x)+b \sin \theta,  \tag{16}\\
\bar{\theta}=\theta+\omega \bmod 2 \pi
\end{array}\right.
$$

where $b>0$ and now $h$ is defined as

$$
h(x)=\left\{\begin{array}{rll}
-\frac{\pi}{2} & \text { if } & x<0 \\
0 & \text { if } & x=0 \\
\frac{\pi}{2} & \text { if } & x>0
\end{array}\right.
$$

As before, we assume $\omega$ to be irrational ( $\omega \notin 2 \pi \mathbb{Q}$ ). We are interested in the existence of invariant curves, that in this case do not need to be continuous.

It follows from (7) that, when $a$ tends to infinity, the critical value $b^{*}(a)$ tends to $\pi / 2$. It is also immediate to see that, if $0<b<\pi / 2$, the map (16) has two smooth attracting invariant curves given by $\varphi_{0}(\theta)=\pi / 2+b \sin (\theta-\omega)$ and $\gamma_{0}(\theta)=-\pi / 2+b \sin (\theta-\omega)$.

### 3.1 The case $b=\pi / 2$

This case is similar to the previous one, in the sense that for a total measure set of values of $\theta$ the attractor lies on two disjoint attracting curves, $\varphi_{0}(\theta)=\pi / 2+(\pi / 2) \sin (\theta-\omega)$ and $\gamma_{0}(\theta)=-\pi / 2+(\pi / 2) \sin (\theta-\omega)$. The difference with the case $b<\pi / 2$ comes from the orbits for which $\varphi_{0}$ or $\gamma_{0}$ takes the value zero. It is clear that $\gamma_{0}(\theta)$ is zero when $\theta=\pi / 2+\omega$ and $\varphi_{0}(\theta)$ is zero when $\theta=3 \pi / 2+\omega$. We distinguish two cases, based on the $\operatorname{sign}$ of $\sin (\pi / 2+\omega)$ (note that this value cannot be zero since then $\omega$ would belong to $2 \pi \mathbb{Q}$ ).

The first case is when $\sin (\pi / 2+\omega)<0$ (and, therefore, $\sin (3 \pi / 2+\omega)>0$ ). In this case there are two disjoint attracting curves. One of these curves coincides with $\gamma_{0}$ except for one point: at $\pi / 2+2 \omega$ it takes the value $(\pi / 2) \sin (\pi / 2+\omega)$. The other curve coincides with $\varphi_{0}$ except for the point $3 \pi / 2+2 \omega$ where it takes the value $(\pi / 2) \sin (3 \pi / 2+\omega)$.

The second case is when $\sin (\pi / 2+\omega)>0$ (and, therefore, $\sin (3 \pi / 2+\omega)<0$ ). Here, the orbit on the curve $\gamma_{0}$ that reaches the value $x=0$ for $\theta=\pi / 2+\omega$ is sent to $x=(\pi / 2) \sin (\pi / 2+\omega)$ which is negative so the next iterate falls on the curve $\varphi_{0}$ and stays there. Similarly, there is an orbit on $\varphi_{0}$ that ends on $\gamma_{0}$. So the attracting set can be described as the union of two discontinuous curves.

### 3.2 The case $b>\pi / 2$

As we will see, in this case there is only one invariant curve, which is attracting and discontinuous.
Theorem 3.1. When $b>\frac{\pi}{2}$ the dynamical system (16) has a unique discontinuous attracting invariant curve. Moreover,
a) the number of discontinuities is at least 2,
b) when $b \rightarrow\left(\frac{\pi}{2}\right)^{+}$the number of discontinuities tends to infinity.

Proof: As before, we denote $\varphi_{0}(\theta)=\pi / 2+b \sin (\theta-\omega)$ and $\gamma_{0}(\theta)=-\pi / 2+b \sin (\theta-\omega)$ that are not invariant since now $b>\pi / 2$. It is clear that the set between the two curves,

$$
A=\left\{(x, \theta) \in \mathbb{R} \times \mathbb{T}^{1} \mid \gamma_{0}(\theta) \leq x \leq \varphi_{0}(\theta)\right\},
$$

is invariant by the map $F_{b}$. It is clear that the set

$$
\Omega=\bigcap_{n=0}^{\infty} F_{b}^{n}(A)
$$

is non-empty and satisfies that $F_{b}(\Omega)=\Omega$. Let us define $\gamma_{n}(\theta)=h\left(\gamma_{n-1}(\theta-\omega)\right)+b \sin (\theta-\omega)$ and $\varphi_{n}(\theta)=h\left(\varphi_{n-1}(\theta-\omega)\right)+b \sin (\theta-\omega)$. Both sequences are bounded, $\left\{\gamma_{n}\right\}_{n}$ is increasing and
$\left\{\varphi_{n}\right\}_{n}$ is decreasing. Therefore, we can define the functions $\gamma_{\infty}$ and $\varphi_{\infty}$ as the pointwise limit of these sequences. Moreover, as $\gamma_{0}$ is the image of $\varphi_{0}$ by the symmetry $S$ (and viceversa), we have that $\gamma_{n}$ is also the image of $\varphi_{n}$ by $S$, and then $\gamma_{\infty}$ is the image of $\varphi_{\infty}$ by $S$, and viceversa.

If we define the map $\sigma$ as $\sigma(\theta)=b \sin (\theta-\omega)$ we have that

$$
\operatorname{Im}\left(\gamma_{\infty}\right) \cup \operatorname{Im}\left(\varphi_{\infty}\right) \subset \operatorname{Im}\left(\gamma_{0}\right) \cup \operatorname{Im}\left(\varphi_{0}\right) \cup \operatorname{Im}(\sigma)
$$

If $\theta$ is such that $\varphi_{\infty}(\theta) \neq 0$, then we can consider $\varphi_{n}(\theta)=h\left(\varphi_{n-1}(\theta-\omega)\right)+b \sin (\theta-\omega)$ and taking limits to both sides we can obtain $\varphi_{\infty}(\theta)=h\left(\varphi_{\infty}(\theta-\omega)\right)+b \sin (\theta-\omega)$. Analogously, if $\gamma_{\infty}(\theta) \neq 0$ we can obtain $\gamma_{\infty}(\theta)=h\left(\gamma_{\infty}(\theta-\omega)\right)+b \sin (\theta-\omega)$. Let us denote by $\theta_{1,2}^{(0)}$ the two zeros of $\varphi_{0}$ and $\theta_{3,4}^{(0)}$ the two zeros of $\gamma_{0}$. Let us consider the set $\mathcal{T}$ defined as the result of removing the set of values $\left\{\left(\theta_{j}^{(0)}+k \omega\right) \bmod 2 \pi \mid k \in \mathbb{N}, 1 \leq j \leq 4\right\}$ from $\mathbb{T}^{1}$. Let us define the map $\tilde{\varphi}_{\infty}$ as follows: if $\theta \in \mathcal{T}, \tilde{\varphi}_{\infty}(\theta)=\varphi_{\infty}(\theta)$ and, if $\theta \notin \mathcal{T}$, we define its value recurrently

$$
\tilde{\varphi}_{\infty}\left(\theta_{j}^{(0)}+k \omega\right)=h\left(\tilde{\varphi}_{\infty}\left(\theta_{j}^{(0)}+(k-1) \omega\right)\right)+b \sin \left(\tilde{\varphi}_{\infty}\left(\theta_{j}^{(0)}+(k-1) \omega\right)\right), \quad k \geq 1
$$

We note that $\tilde{\varphi}_{\infty}$ is an invariant, not necessarily continuous, curve. The curve $\tilde{\gamma}_{\infty}$ can be defined analogously. On the other hand, there exist an open interval $I_{1}$ on which the two functions $\varphi_{0}$ and $\gamma_{0}$ are strictly positive (and hence, due to the symmetry $S$, they are strictly negative on the interval $I_{1}+\pi$ ). This implies that $\varphi_{1}$ and $\gamma_{1}$ coincide on $I_{1}$. As

$$
\gamma_{n}(\theta) \leq \tilde{\gamma}_{\infty}(\theta) \leq \tilde{\varphi}_{\infty}(\theta) \leq \varphi_{n}(\theta), \quad \theta \in \mathbb{T}^{1}, \quad n \geq 0,
$$

then $\tilde{\gamma}_{\infty}$ and $\tilde{\varphi}_{\infty}$ coincide on the interval $I_{1}$. As both are invariant, they must coincide. Finally, let us note that $\tilde{\gamma}_{\infty}$ and $\gamma_{\infty}$ (respectively $\tilde{\varphi}_{\infty}$ and $\varphi_{\infty}$ ) are identical except on a finite (and strictly positive) number of values of $\theta$. This implies that $\gamma_{\infty}$ and $\varphi_{\infty}$ are also identical except on a finite number of points. In summary, we have shown that there exists an attracting invariant curve $\tilde{\varphi}_{\infty}$ (or $\tilde{\gamma}_{\infty}$ since they are identical) which is unique, and that the number of discontinuities is at least 2 , so item a) is proved.

To prove item b), let us define $\theta_{0}$ as the minimum of $\varphi_{0}$, so $\theta_{0}=-\pi / 2+\omega$. The symmetry $S$ implies that $\theta_{1}=\pi / 2+\omega$ is the maximum of $\gamma_{0}$. Let us also define the sets

$$
\Theta_{M}^{(i)}=\left\{\theta_{i}+k \omega(\bmod 2 \pi) \mid k=1, \ldots, M\right\}, \quad i=0,1 .
$$

Note that, for all $M \in \mathbb{N}$, the fact $\omega \notin 2 \pi \mathbb{Q}$ implies $\Theta_{M}^{(0)} \cap \Theta_{M}^{(1)}=\emptyset$ and $0 \notin \Theta_{M}^{(0)} \cup \Theta_{M}^{(1)}$ (if not, then $\omega \in 2 \pi \mathbb{Q})$. Let us see that, for a given number of discontinuities $n_{d}$, there exist a value $b_{n_{d}}>\pi / 2$ such that the attracting curve has at least $n_{d}$ discontinuities for $b \in\left(\pi / 2, b_{n_{d}}\right]$. To this end, we will construct $n_{d}$ intervals where the attractor is on $\varphi_{0}$ and another $n_{d}$ intervals, interspersed with the previous intervals, where the attractor is on $\gamma_{0}$. This will guarantee the existence of more than $n_{d}$ discontinuities. To this end, let us choose a value $M_{n_{d}}$ such that each of the sets $\Theta_{M}^{(0)}$ and $\Theta_{M}^{(1)}$ contain $n_{d}$ points that are interspersed with points of the other set: in other words, there exist angles $0<\alpha_{1}<\cdots<\alpha_{n_{d}}<2 \pi$ in $\Theta_{M}^{(0)}$ and $0<\beta_{1}<\cdots<\beta_{n_{d}}<2 \pi$ in $\Theta_{M}^{(1)}$ such that $\alpha_{j}<\beta_{j}$ for $j=1, \ldots, n_{d}$. The existence of these angles is ensured by the fact that both $\Theta_{M}^{(0)}$ and $\Theta_{M}^{(1)}$ fill densely the circle when $M \rightarrow \infty$. The final step is to construct small open intervals, $I_{\alpha_{j}}, I_{\beta_{j}}$, with $\alpha_{j} \in I_{\alpha_{j}}$ and $\beta_{j} \in I_{\beta_{j}}$ and such that the attractor is on $\gamma_{0}$ for $I_{\alpha_{j}}$ and on $\varphi_{0}$ for $I_{\beta_{j}}$. Let us start by defining the intervals $I_{\theta_{i}}=\left(\theta_{i}-\delta, \theta_{i}+\delta\right), i=0,1$. Let $\delta>0$
be such that all the intervals $I_{\theta_{i}+k \omega}\left(i=0,1,0 \leq k \leq M_{n_{d}}\right)$ are disjoint. We note that $\delta$ goes to zero when $n_{d}$ goes to infinity. Now, let us choose a value of $b, b_{n_{d}}$, such that the interval centered in $\theta_{0}$ and of radius $\delta, I_{\theta_{0}}$, is the largest open interval centered in $\theta_{0}$ such that $\varphi_{0}\left(I_{\theta_{0}}\right)<0$. By symmetry, $I_{\theta_{1}}=I_{\theta_{0}}+\pi$ is the largest open interval centered in $\theta_{1}$ such that $\gamma_{0}\left(I_{\theta_{1}}\right)>0$. Then, the attracting invariant curve $\tilde{\varphi}_{\infty}$ satisfies that $\tilde{\varphi}_{\infty}(\theta)<0$ if $\theta \in I_{\theta_{0}}+k \omega$ and $\tilde{\varphi}_{\infty}(\theta)>0$ if $\theta \in I_{\theta_{1}}+k \omega, 1 \leq k \leq M_{n_{d}}$. In particular, as $n_{d}$ of these intervals are interspersed, we ensure the existence of $n_{d}$ discontinuities. We note that this number of discontinuities is ensured for any value of $b$ such that $\pi / 2<b<b_{n_{d}}$.

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